



# **ELECTROSTATIC**

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# INTRODUCTION

- **Electrostatic** is fundamental concepts that are applicable to static(or time-invariant) electric fields in free space(or vacuum).An electrostatic field is produced by a static charge distribution.
- Study of electrostatics by investigating the two fundamental laws governing electrostatic fields:
  - Coulomb's law, and
  - Gauss'slaw.



# COULOMB'S LAW

**Coulomb's law** states that the force  $F$  between two point charges  $Q_1$  and  $Q_2$  is:

1. Along the line joining them
2. Directly proportional to the product  $Q_1Q_2$  of the charges
3. Inversely proportional to the square of the distance  $R$  between them.<sup>3</sup>

Expressed mathematically,

$$F = \frac{k Q_1 Q_2}{R^2}$$

Where  $k$  is the proportionality constant. In SI units, charges are in coulombs (C), The distance  $R$  is in meters(m), and the force  $F$  is in newton (N) so that

$$\begin{aligned} \epsilon_0 &= 8.854 \times 10^{-12} \approx \frac{10^{-9}}{36\pi} \text{ F/m} \\ \text{or } k &= \frac{1}{4\pi\epsilon_0} \approx 9 \times 10^9 \text{ m/F} \end{aligned}$$



# VECTOR FORM OF COULOMB'S LAW

$$\mathbf{F}_{12} = \frac{Q_1 Q_2}{4\pi\epsilon_0 R^2} \mathbf{a}_{R_{12}}$$

where

$$\mathbf{R}_{12} = \mathbf{r}_2 - \mathbf{r}_1$$

$$R = |\mathbf{R}_{12}|$$

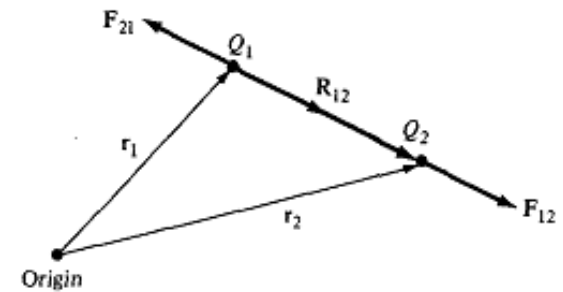
$$\mathbf{a}_{R_{12}} = \frac{\mathbf{R}_{12}}{R}$$

therefore

$$\mathbf{F}_{12} = \frac{Q_1 Q_2}{4\pi\epsilon_0 R^3} \mathbf{R}_{12}$$

or

$$\mathbf{F}_{12} = \frac{Q_1 Q_2 (\mathbf{r}_2 - \mathbf{r}_1)}{4\pi\epsilon_0 |\mathbf{r}_2 - \mathbf{r}_1|^3}$$



# COULOMB'S LAW

(IMP. PTS TO BE NOTED)

1.  $\mathbf{F}_{21} = -\mathbf{F}_{12}$
2. Like charges (charges of the same sign) repel each other while unlike charges attract. This is illustrated in Figure 4.2.
3. The distance  $R$  between the charged bodies  $Q_1$  and  $Q_2$  must be large compared with the linear dimensions of the bodies; that is,  $Q_1$  and  $Q_2$  must be point charges.
4.  $Q_1$  and  $Q_2$  must be static (at rest).
5. The signs of  $Q_1$  and  $Q_2$  must be taken into account



# PRINCIPLE OF SUPERPOSITION

If we have more than two point charges, we can use the *principle of superposition* to determine the force on a particular charge. The principle states that if there are  $N$  charges  $Q_1, Q_2, \dots, Q_N$  located, respectively, at points with position vectors  $\mathbf{r}_1, \mathbf{r}_2, \dots, \mathbf{r}_N$ , the resultant force  $\mathbf{F}$  on a charge  $Q$  located at point  $\mathbf{r}$  is the vector sum of the forces exerted on  $Q$  by each of the charges  $Q_1, Q_2, \dots, Q_N$ . Hence:

$$\mathbf{F} = \frac{QQ_1(\mathbf{r} - \mathbf{r}_1)}{4\pi\epsilon_0|\mathbf{r} - \mathbf{r}_1|^3} + \frac{QQ_2(\mathbf{r} - \mathbf{r}_2)}{4\pi\epsilon_0|\mathbf{r} - \mathbf{r}_2|^3} + \dots + \frac{QQ_N(\mathbf{r} - \mathbf{r}_N)}{4\pi\epsilon_0|\mathbf{r} - \mathbf{r}_N|^3}$$

OR

$$\mathbf{F} = \frac{Q}{4\pi\epsilon_0} \sum_{k=1}^N \frac{Q_k(\mathbf{r} - \mathbf{r}_k)}{|\mathbf{r} - \mathbf{r}_k|^3}$$



# ELECTRIC FIELD INTENSITY

The **electric field intensity** (or **electric field strength**) **E** is the force per unit charge when placed in the electric field.

Thus

$$\mathbf{E} = \lim_{Q \rightarrow 0} \frac{\mathbf{F}}{Q}$$

or simply

$$\mathbf{E} = \frac{\mathbf{F}}{Q}$$

The electric field intensity **E** is obviously in the direction of the force **F** and is measured in newton/coulomb or volts/meter.





# ELECTRIC FIELD INTENSITY

- The electric field intensity at point  $\mathbf{r}$  due to a point charge located at  $\mathbf{r}'$  is

$$\mathbf{E} = \frac{Q}{4\pi\epsilon_0 R^2} \mathbf{a}_R = \frac{Q(\mathbf{r} - \mathbf{r}')}{4\pi\epsilon_0 |\mathbf{r} - \mathbf{r}'|^3}$$

For  $N$  point charges  $Q_1, Q_2, \dots, Q_N$  located at  $\mathbf{r}_1, \mathbf{r}_2, \dots, \mathbf{r}_N$ , the electric field intensity at point  $\mathbf{r}$  is obtained from eqs. (4.8) and (4.10) as

$$\mathbf{E} = \frac{Q_1(\mathbf{r} - \mathbf{r}_1)}{4\pi\epsilon_0 |\mathbf{r} - \mathbf{r}_1|^3} + \frac{Q_2(\mathbf{r} - \mathbf{r}_2)}{4\pi\epsilon_0 |\mathbf{r} - \mathbf{r}_2|^3} + \dots + \frac{Q_N(\mathbf{r} - \mathbf{r}_N)}{4\pi\epsilon_0 |\mathbf{r} - \mathbf{r}_N|^3}$$

or

$$\mathbf{E} = \frac{1}{4\pi\epsilon_0} \sum_{k=1}^N \frac{Q_k(\mathbf{r} - \mathbf{r}_k)}{|\mathbf{r} - \mathbf{r}_k|^3}$$





## EXAMPLE 1

Point charges 1 mC and  $-2$  mC are located at  $(3, 2, -1)$  and  $(-1, -1, 4)$ , respectively. Calculate the electric force on a 10-nC charge located at  $(0, 3, 1)$  and the electric field intensity at that point.

**Solution:**

$$\begin{aligned} \mathbf{F} &= \sum_{k=1,2} \frac{QQ_k}{4\pi\epsilon_0 R^2} \mathbf{a}_R = \sum_{k=1,2} \frac{QQ_k(\mathbf{r} - \mathbf{r}_k)}{4\pi\epsilon_0 |\mathbf{r} - \mathbf{r}_k|^3} \\ &= \frac{Q}{4\pi\epsilon_0} \left\{ \frac{10^{-3}[(0, 3, 1) - (3, 2, -1)]}{|(0, 3, 1) - (3, 2, -1)|^3} - \frac{2 \cdot 10^{-3}[(0, 3, 1) - (-1, -1, 4)]}{|(0, 3, 1) - (-1, -1, 4)|^3} \right\} \\ &= \frac{10^{-3} \cdot 10 \cdot 10^{-9}}{4\pi \cdot \frac{10^{-9}}{36\pi}} \left[ \frac{(-3, 1, 2)}{(9 + 1 + 4)^{3/2}} - \frac{2(1, 4, -3)}{(1 + 16 + 9)^{3/2}} \right] \\ &= 9 \cdot 10^{-2} \left[ \frac{(-3, 1, 2)}{14\sqrt{14}} + \frac{(-2, -8, 6)}{26\sqrt{26}} \right] \\ \mathbf{F} &= -6.507\mathbf{a}_x - 3.817\mathbf{a}_y + 7.506\mathbf{a}_z \text{ mN} \end{aligned}$$



## EXAMPLE 1

At that point,

$$\begin{aligned}\mathbf{E} &= \frac{\mathbf{F}}{Q} \\ &= (-6.507, -3.817, 7.506) \cdot \frac{10^{-3}}{10 \cdot 10^{-9}} \\ \mathbf{E} &= -650.7\mathbf{a}_x - 381.7\mathbf{a}_y + 750.6\mathbf{a}_z \text{ kV/m}\end{aligned}$$



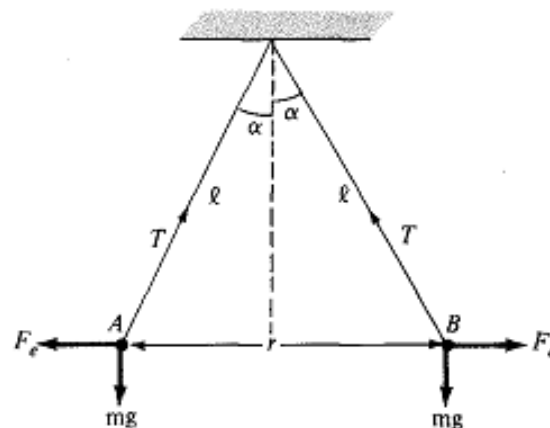
## EXAMPLE 2

Two point charges of equal mass  $m$ , charge  $Q$  are suspended at a common point by two threads of negligible mass and length  $\ell$ . Show that at equilibrium the inclination angle  $\alpha$  of each thread to the vertical is given by

$$Q^2 = 16\pi \epsilon_0 mg\ell^2 \sin^2 \alpha \tan \alpha$$

If  $\alpha$  is very small, show that

$$\alpha = \sqrt[3]{\frac{Q^2}{16\pi\epsilon_0 mg\ell^2}}$$



## EXAMPLE 2

### Solution:

Consider the system of charges as shown in Figure 4.3 where  $F_e$  is the electric or coulomb force,  $T$  is the tension in each thread, and  $mg$  is the weight of each charge. At  $A$  or  $B$

$$T \sin \alpha = F_e$$

$$T \cos \alpha = mg$$

Hence,

$$\frac{\sin \alpha}{\cos \alpha} = \frac{F_e}{mg} = \frac{1}{mg} \cdot \frac{Q^2}{4\pi\epsilon_0 r^2}$$

But

$$r = 2\ell \sin \alpha$$

Hence,

$$Q^2 \cos \alpha = 16\pi\epsilon_0 mg\ell^2 \sin^3 \alpha$$

or

$$Q^2 = 16\pi\epsilon_0 mg\ell^2 \sin^2 \alpha \tan \alpha$$

as required. When  $\alpha$  is very small

$$\tan \alpha \approx \alpha \approx \sin \alpha$$



## EXAMPLE 2

and so

$$Q^2 = 16\pi\epsilon_0 m g \ell^2 \alpha^3$$

or

$$\alpha = \sqrt[3]{\frac{Q^2}{16\pi\epsilon_0 m g \ell^2}}$$

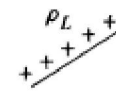


# ELECTRIC FIELD DUE TO CONTINUOUS CHARGE DISTRIBUTION

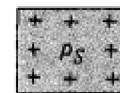
- So far we have only considered forces and electric fields due to point charges, which are essentially charges occupying very small physical space.
- It is also possible to have continuous charge distribution along a line, on a surface, or in a volume.
  - Line charge density by  $\rho_L$  (in C/m),
  - surface charge density by  $\rho_S$  (in C/m<sup>2</sup>), and
  - volume charge density and  $\rho_V$  (in C/m<sup>3</sup>),



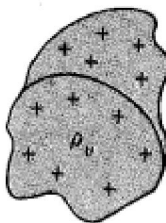
Point charge



Line charge



Surface charge



Volume charge

# ELECTRIC FIELD DUE TO CONTINUOUS CHARGE DISTRIBUTION

- The charge element  $dQ$  and total charge  $Q$  due to these continuous charge distribution are

$$dQ = \rho_L dl \rightarrow Q = \int_L \rho_L dl \quad (\text{line charge})$$

$$dQ = \rho_S dS \rightarrow Q = \int_S \rho_S dS \quad (\text{surface charge})$$

$$dQ = \rho_V dv \rightarrow Q = \int_V \rho_V dv \quad (\text{volume charge})$$





# ELECTRIC FIELD DUE TO CONTINUOUS CHARGE DISTRIBUTION

- Electric field due to this charge distribution will be

$$\mathbf{E} = \int \frac{\rho_L dl}{4\pi\epsilon_0 R^2} \mathbf{a}_R \quad (\text{line charge})$$

$$\mathbf{E} = \int \frac{\rho_S dS}{4\pi\epsilon_0 R^2} \mathbf{a}_R \quad (\text{surface charge})$$

$$\mathbf{E} = \int \frac{\rho_V dv}{4\pi\epsilon_0 R^2} \mathbf{a}_R \quad (\text{volume charge})$$



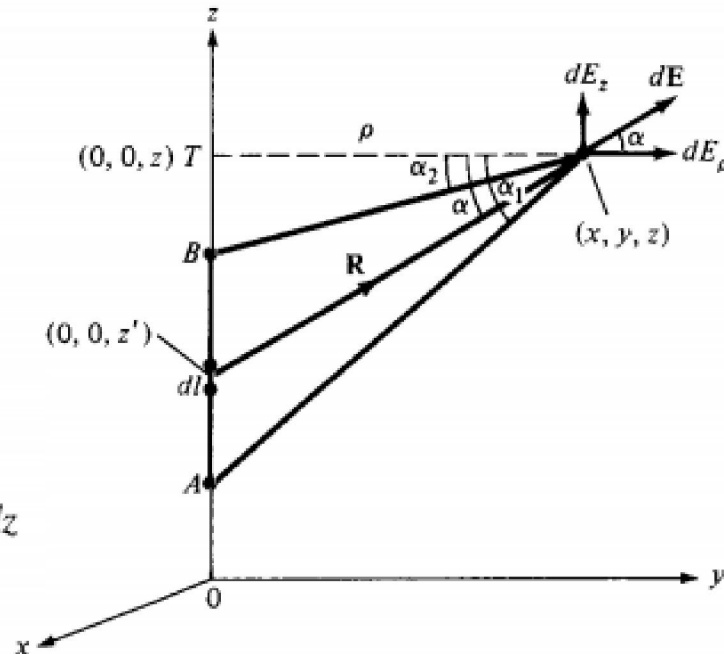
# A LINE CHARGE

- Consider a line charge with uniform charge density  $\rho_L$  extending from  $A$  to  $B$  along the  $z$ -axis as shown in Figure.
- The charge element  $dQ$  associated with element  $dl = dz$  of the line is

$$dQ = \rho_L dl = \rho_L dz$$

and hence the total charge  $Q$  is

$$Q = \int_{z_A}^{z_B} \rho_L dz$$



# A LINE CHARGE

$$\mathbf{E} = \int \frac{\rho_L dl}{4\pi\epsilon_0 R^2} \mathbf{a}_R \quad (\text{line charge})$$

$$dl = dz'$$

$$\mathbf{R} = (x, y, z) - (0, 0, z') = x\mathbf{a}_x + y\mathbf{a}_y + (z - z')\mathbf{a}_z$$

or

$$\mathbf{R} = \rho\mathbf{a}_\rho + (z - z')\mathbf{a}_z$$

$$R^2 = |\mathbf{R}|^2 = x^2 + y^2 + (z - z')^2 = \rho^2 + (z - z')^2$$

$$\frac{\mathbf{a}_R}{R^2} = \frac{\mathbf{R}}{|\mathbf{R}|^3} = \frac{\rho\mathbf{a}_\rho + (z - z')\mathbf{a}_z}{[\rho^2 + (z - z')^2]^{3/2}}$$



# A LINE CHARGE

$$\mathbf{E} = \frac{\rho_L}{4\pi\epsilon_0} \int \frac{\rho \mathbf{a}_\rho + (z - z') \mathbf{a}_z}{[\rho^2 + (z - z')^2]^{3/2}} dz'$$

To evaluate this, it is convenient that we define  $\alpha$ ,  $\alpha_1$ , and  $\alpha_2$  as in Figure

$$R = [\rho^2 + (z - z')^2]^{1/2} = \rho \sec \alpha$$

$$z' = OT - \rho \tan \alpha, \quad dz' = -\rho \sec^2 \alpha d\alpha$$

Hence, eq. becomes

$$\begin{aligned} \mathbf{E} &= \frac{-\rho_L}{4\pi\epsilon_0} \int_{\alpha_1}^{\alpha_2} \frac{\rho \sec^2 \alpha [\cos \alpha \mathbf{a}_\rho + \sin \alpha \mathbf{a}_z] d\alpha}{\rho^2 \sec^2 \alpha} \\ &= -\frac{\rho_L}{4\pi\epsilon_0 \rho} \int_{\alpha_1}^{\alpha_2} [\cos \alpha \mathbf{a}_\rho + \sin \alpha \mathbf{a}_z] d\alpha \end{aligned}$$

Thus for a *finite line charge*,

$$\mathbf{E} = \frac{\rho_L}{4\pi\epsilon_0 \rho} [-(\sin \alpha_2 - \sin \alpha_1) \mathbf{a}_\rho + (\cos \alpha_2 - \cos \alpha_1) \mathbf{a}_z]$$



# A LINE CHARGE

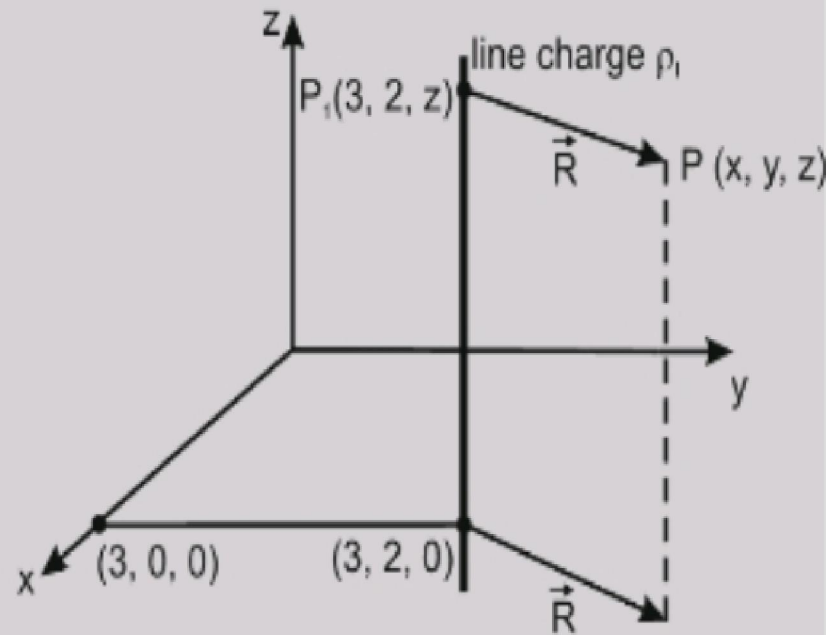
As a special case, for an *infinite line charge*, point  $B$  is at  $(0, 0, \infty)$  and  $A$  at  $(0, 0, -\infty)$  so that  $\alpha_1 = \pi/2$ ,  $\alpha_2 = -\pi/2$ ; the  $z$ -component vanishes and eq. (4.20) becomes

$$\mathbf{E} = \frac{\rho_L}{2\pi\epsilon_0\rho} \mathbf{a}_\rho$$



# EXAMPLE

An infinite uniform line charge is placed parallel to the  $z$ -axis at  $x = 3$  and  $y = 2$ . A point  $P(x, y, z)$  is located at a distance  $R$  from the line charge as shown in the following figure. Find the electric field intensity  $\vec{E}$  at point  $P$ .



# EXAMPLE

**Solution** We know that the electric field intensity  $\vec{E}$  for the infinite line charge along  $z$ -axis is:

$$\vec{E} = \frac{\rho_l}{2\pi\epsilon_0\rho} \hat{a}_\rho$$

Replacing  $\rho$  in the above equation by the radial distance  $R$  between the line charge and the point  $P$  and assuming that  $\hat{a}_\rho = \hat{a}_R$ , we can write from the figure that:

$$\vec{R} = (x, y, z) - (3, 2, z) = (x - 3) \hat{a}_x + (y - 2) \hat{a}_y$$

$$|\vec{R}| = \sqrt{(x - 3)^2 + (y - 2)^2}$$

$$\hat{a}_R = \frac{\vec{R}}{|\vec{R}|} = \frac{(x - 3) \hat{a}_x + (y - 2) \hat{a}_y}{\sqrt{(x - 3)^2 + (y - 2)^2}}$$

Thus,

$$\vec{E} = \frac{\rho_l}{2\pi\epsilon_0} \frac{(x - 3) \hat{a}_x + (y - 2) \hat{a}_y}{(x - 3)^2 + (y - 2)^2}$$





## EXAMPLE

A uniform line charge density of 10 nC/m is located at  $x = 0, y = 2$ , while another uniform line charge density of  $-10$  nC/m is located at  $x = 0, y = -2$ . Determine the electric field at the origin.

**Solution** As we are given two infinitely long charged lines, the total electric field  $\vec{E}$  at any point is the vector sum of the electric field at that point due to individual line charges. That is,

$$\vec{E} = \frac{\rho_{l_1}}{2\pi\epsilon_0\rho_1} \hat{a}_{\rho_{l_1}} + \frac{\rho_{l_2}}{2\pi\epsilon_0\rho_2} \hat{a}_{\rho_{l_2}}$$

Given that  $\rho_{l_1} = 10$  nC/m and  $\rho_{l_2} = -10$  nC/m. Also, the location of  $\rho_{l_1}$  is  $x = 0, y = 2$  and the location of  $\rho_{l_2}$  is  $x = 0, y = -2$ , and we have to evaluate electric field at the origin, that is, at  $(0, 0, 0)$ . Therefore:



## EXAMPLE

$$\rho_1 = \sqrt{0^2 + 2^2} = 2 \quad \text{and} \quad \rho_2 = \sqrt{0^2 + (-2)^2} = 2$$

$$\hat{a}_{\rho_1} = \frac{\bar{\rho}_{l_1}}{\rho_1} = \frac{2\hat{a}_y}{2} = \hat{a}_y \quad \text{and} \quad \hat{a}_{\rho_2} = \frac{\rho_{l_2}}{\rho_2} = \frac{-2\hat{a}_y}{2} = -\hat{a}_y$$

As we know that  $\frac{1}{4\pi\epsilon_0} = 9 \times 10^9 \text{ m/F}$ . This implies that  $\frac{1}{2\pi\epsilon_0} = 18 \times 10^9 \text{ m/F}$ .

Thus,  $\vec{E}$  at the origin is computed as follows.

$$\vec{E} = \frac{10 \times 10^{-9} \times 18 \times 10^9}{2} \hat{a}_y + \frac{-10 \times 10^{-9} \times 18 \times 10^9}{2} (-\hat{a}_y) = 180\hat{a}_y \text{ V/m}$$



# SURFACE CHARGE

- Consider an infinite sheet of charge in the  $xy$ -plane with uniform charge density  $\rho_s$ . The charge associated with an elemental area  $dS$  is

$$dQ = \rho_s dS$$

and hence the total charge is

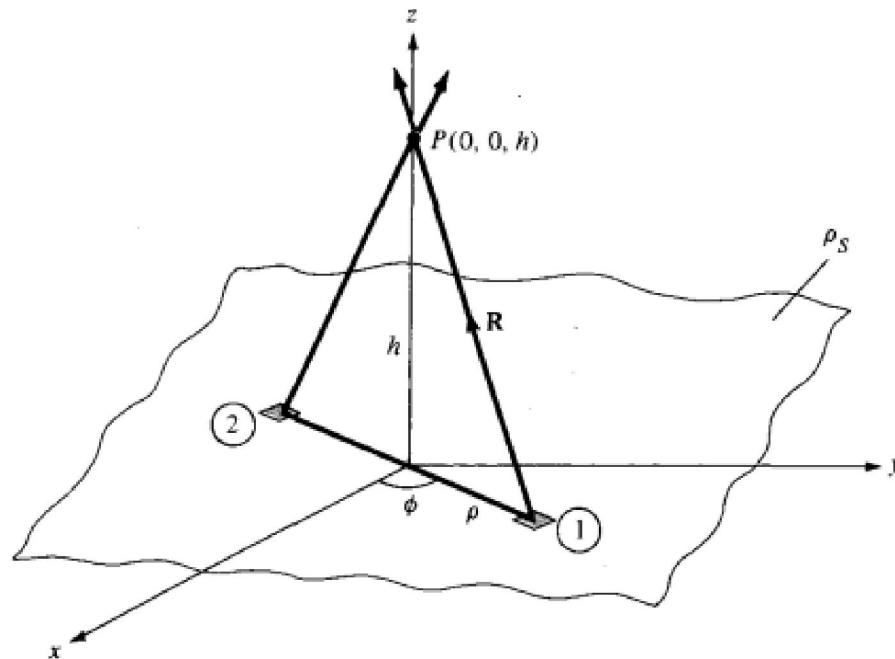
$$Q = \int \rho_s dS$$



# SURFACE CHARGE

- The contribution to the E field at point  $P(0, 0, h)$  by the elemental surface 1 shown in Figure

$$d\mathbf{E} = \frac{dQ}{4\pi\epsilon_0 R^2} \mathbf{a}_R$$



# SURFACE CHARGE

$$\mathbf{R} = \rho(-\mathbf{a}_\rho) + h\mathbf{a}_z, \quad R = |\mathbf{R}| = [\rho^2 + h^2]^{1/2}$$

$$\mathbf{a}_R = \frac{\mathbf{R}}{R}, \quad dQ = \rho_S dS = \rho_S \rho d\phi d\rho$$

Substitution of these terms gives

$$d\mathbf{E} = \frac{\rho_S \rho d\phi d\rho [-\rho\mathbf{a}_\rho + h\mathbf{a}_z]}{4\pi\epsilon_0[\rho^2 + h^2]^{3/2}}$$

Due to the symmetry of the charge distribution, for every element 1, there is a corresponding element 2 whose contribution along  $\mathbf{a}_\rho$  cancels that of element 1, as illustrated in Figure . Thus the contributions to  $E_\rho$  add up to zero so that  $\mathbf{E}$  has only z-component.



# SURFACE CHARGE

$$\begin{aligned}\mathbf{E} &= \int d\mathbf{E}_z = \frac{\rho_S}{4\pi\epsilon_0} \int_{\phi=0}^{2\pi} \int_{\rho=0}^{\infty} \frac{h\rho d\rho d\phi}{[\rho^2 + h^2]^{3/2}} \mathbf{a}_z \\ &= \frac{\rho_S h}{4\pi\epsilon_0} 2\pi \int_0^{\infty} [\rho^2 + h^2]^{-3/2} \frac{1}{2} d(\rho^2) \mathbf{a}_z \\ &= \frac{\rho_S h}{2\epsilon_0} \left\{ -[\rho^2 + h^2]^{-1/2} \right\}_0^{\infty} \mathbf{a}_z \\ \mathbf{E} &= \frac{\rho_S}{2\epsilon_0} \mathbf{a}_z\end{aligned}$$

that is,  $\mathbf{E}$  has only  $z$ -component if the charge is in the  $xy$ -plane. In general, for an *infinite sheet* of charge

$$\mathbf{E} = \frac{\rho_S}{2\epsilon_0} \mathbf{a}_n$$



# SURFACE CHARGE

- In a parallel plate capacitor, the electric field existing between the two plates having equal and opposite charges is given by

$$\mathbf{E} = \frac{\rho_S}{2\epsilon_0} \mathbf{a}_n + \frac{-\rho_S}{2\epsilon_0} (-\mathbf{a}_n) = \frac{\rho_S}{\epsilon_0} \mathbf{a}_n$$





# EXAMPLE

An infinite sheet in  $xy$ -plane extending from  $-\infty$  to  $\infty$  in both directions has a uniform charge density of  $5 \text{ nC/m}^2$ . Find the electric field at  $z = 10 \text{ cm}$ .

**Solution:** Given that  $\rho_s = 5 \text{ nC/m}^2 = 5 \times 10^{-9} \text{ C/m}^2$

As the given sheet lies in  $xy$ -plane and the field point is on  $z$ -axis, the electric field  $\vec{E}$  can be given as:

$$\vec{E} = \frac{\rho_s}{2\epsilon_0} \hat{a}_z$$

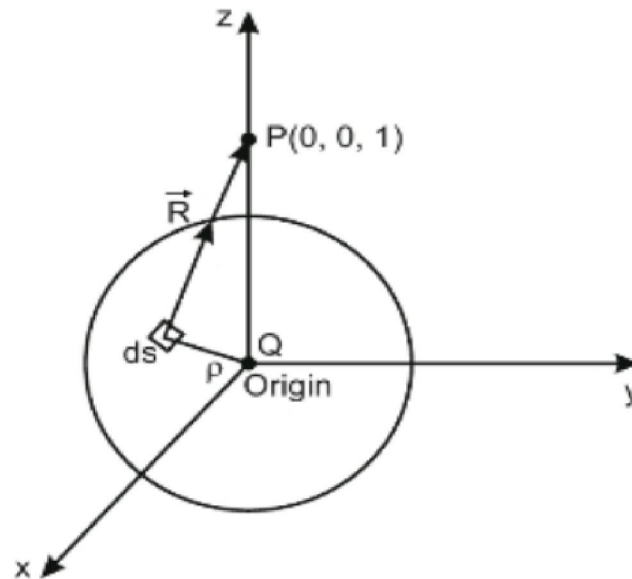
$$\Rightarrow \vec{E} = \frac{5 \times 10^{-9}}{2 \times 8.854 \times 10^{-12}} \hat{a}_z = 282.3 \hat{a}_z \text{ V/m}$$



# EXAMPLE

A disk located at  $0 < \rho < 1$ ,  $z = 1$  has a uniform charge distribution  $\rho_s = 200 \text{ pC/m}^2$ . A  $30 \text{ }\mu\text{C}$  point charge  $Q$  is located at  $(0, 0, 0)$ . Find the force on  $Q$  due to electric field produced by the disk.

**Solution** Consider small elemental area  $ds$  of the finite uniformly charged disk as shown in the following figure



If the field point  $P$  on  $z$ -axis ( $z = 1$ ) is at a distance  $R$  from  $ds$ , then observe from the figure that:



# EXAMPLE

$$\vec{R} = -\rho\hat{a}_\rho + 1\hat{a}_z = -\rho\hat{a}_\rho + \hat{a}_z$$

⇒

$$|\vec{R}| = \sqrt{(-\rho)^2 + 1} = \sqrt{1 + \rho^2}$$

the differential surface in cylindrical coordinate system along z-direction is:

$$ds = \rho d\phi d\rho$$

Now, the electric field at a point due to sheet charge distribution can be obtained as:

$$\vec{E} = \int_s \frac{\rho_s ds}{4\pi\epsilon_o R} \hat{a}_R = \int_s \frac{\rho_s ds \vec{R}}{4\pi\epsilon_o R^3}$$

On putting values of  $ds$ ,  $R$ , and  $\vec{R}$ , we have:

$$\vec{E} = \int_s \frac{\rho_s (-\rho\hat{a}_\rho + \hat{a}_z) \rho d\phi d\rho}{4\pi\epsilon_o (1 + \rho^2)^{3/2}}$$



# EXAMPLE

Due to symmetry, the electric field along  $\rho$ -direction contributes to zero and we have electric field only along the  $z$ -direction. Thus, ignoring the  $\hat{a}_\rho$  component in the above equation and taking only  $\hat{a}_z$  component, we get  $E_z$  as shown below.

$$\begin{aligned} E_z &= \frac{\rho_s}{4\pi\epsilon_o} \int_s \frac{\rho d\phi d\rho}{(1+\rho^2)^{3/2}} \\ &= \frac{\rho_s}{4\pi\epsilon_o} \int_{\phi=0}^{2\pi} d\phi \int_{\rho=0}^1 \frac{\rho}{(1+\rho^2)^{3/2}} d\rho = \frac{\rho_s}{4\pi\epsilon_o} [\phi]_0^{2\pi} \int_{\rho=0}^1 \frac{\rho}{(1+\rho^2)^{3/2}} d\rho \quad [\text{Given } 0 < \rho < 1] \end{aligned}$$

To solve the integral, let us take  $1 + \rho^2 = v$ . This gives  $\frac{dv}{d\rho} = 2\rho \Rightarrow dv = 2\rho d\rho$ . Thus,  $E_z$  becomes:

$$E_z = \frac{\rho_s}{4\pi\epsilon_o} (2\pi) \frac{1}{2} \int v^{-3/2} dv = -\frac{\rho_s}{4\epsilon_o} [2v^{-1/2}] = -\frac{\rho_s}{2\epsilon_o} [v^{-1/2}]$$

Replacing  $v$  with  $(1 + \rho^2)$  in the above equation and using limits  $0 < \rho < 1$ , we get:

$$E_z = -\frac{\rho_s}{2\epsilon_o} \left[ \frac{1}{\sqrt{1+\rho^2}} \right]_0^1 = \frac{\rho_s}{2\epsilon_o} \left( 1 - \frac{1}{\sqrt{2}} \right)$$

Hence,

$$\vec{E} = \frac{\rho_s}{2\epsilon_o} \left( 1 - \frac{1}{\sqrt{2}} \right) \hat{a}_z$$



## EXAMPLE

Now,  $\vec{F} = Q\vec{E}$ . Given that  $Q = 30 \mu\text{C}$  and  $\rho_s = 200 \text{ pC/m}^2$ . Therefore:

$$\vec{F} = \frac{30 \times 10^{-6} \times 200 \times 10^{-12}}{2 \times \frac{10^{-9}}{36\pi}} \left(1 - \frac{1}{\sqrt{2}}\right) \hat{a}_z = 99.24 \hat{a}_z \mu\text{N}$$



# EXAMPLE

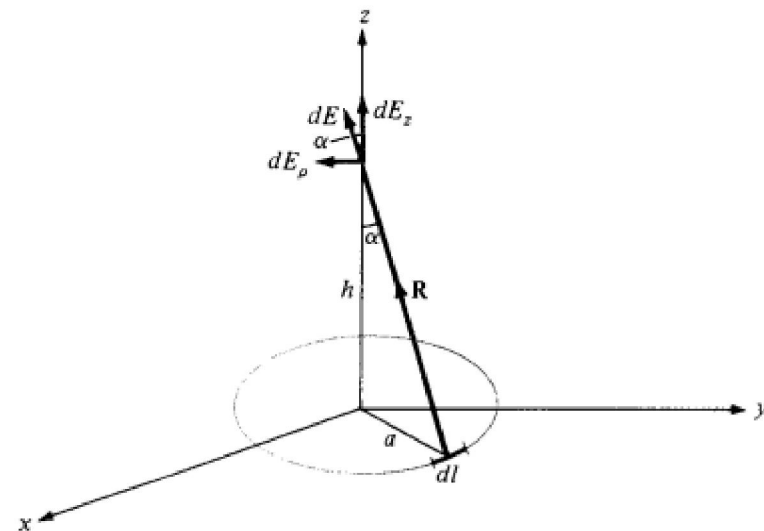
A circular ring of radius  $a$  carries a uniform charge  $\rho_L$  C/m and is placed on the  $xy$ -plane with axis the same as the  $z$ -axis.

(a) Show that

$$\mathbf{E}(0, 0, h) = \frac{\rho_L a h}{2\epsilon_0 [h^2 + a^2]^{3/2}} \mathbf{a}_z$$

(b) What values of  $h$  gives the maximum value of  $\mathbf{E}$ ?

(c) If the total charge on the ring is  $Q$ , find  $\mathbf{E}$  as  $a \rightarrow 0$ .



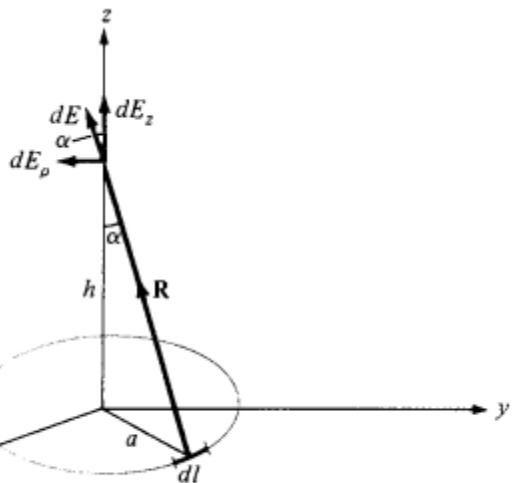
# SOLUTION

$$dl = a d\phi, \quad \mathbf{R} = a(-\mathbf{a}_\rho) + h\mathbf{a}_z$$

$$R = |\mathbf{R}| = [a^2 + h^2]^{1/2}, \quad \mathbf{a}_R = \frac{\mathbf{R}}{R}$$

or

$$\frac{\mathbf{a}_R}{R^2} = \frac{\mathbf{R}}{|\mathbf{R}|^3} = \frac{-a\mathbf{a}_\rho + h\mathbf{a}_z}{[a^2 + h^2]^{3/2}}$$



$$\mathbf{E} = \frac{\rho_L}{4\pi\epsilon_0} \int_{\phi=0}^{2\pi} \frac{(-a\mathbf{a}_\rho + h\mathbf{a}_z)}{[a^2 + h^2]^{3/2}} a d\phi$$





## SOLUTION

- By symmetry radial components are cancelled out and field is only due to Z axis

$$\mathbf{E} = \frac{\rho_L a h \mathbf{a}_z}{4\pi\epsilon_0 [h^2 + a^2]^{3/2}} \int_0^{2\pi} d\phi = \frac{\rho_L a h \mathbf{a}_z}{2\epsilon_0 [h^2 + a^2]^{3/2}}$$

(b)

$$\frac{d|\mathbf{E}|}{dh} = \frac{\rho_L a}{2\epsilon_0} \left\{ \frac{[h^2 + a^2]^{3/2}(1) - \frac{3}{2}(h)2h[h^2 + a^2]^{1/2}}{[h^2 + a^2]^3} \right\}$$

For maximum  $\mathbf{E}$ ,  $\frac{d|\mathbf{E}|}{dh} = 0$ , which implies that

$$[h^2 + a^2]^{1/2} [h^2 + a^2 - 3h^2] = 0$$

$$a^2 - 2h^2 = 0 \quad \text{or} \quad h = \pm \frac{a}{\sqrt{2}}$$



## SOLUTION

(c) Since the charge is uniformly distributed, the line charge density is

$$\rho_L = \frac{Q}{2\pi a}$$

so that

$$\mathbf{E} = \frac{Qh}{4\pi\epsilon_0[h^2 + a^2]^{3/2}} \mathbf{a}_z$$

As  $a \rightarrow 0$

$$\mathbf{E} = \frac{Q}{4\pi\epsilon_0 h^2} \mathbf{a}_z$$

or in general

$$\mathbf{E} = \frac{Q}{4\pi\epsilon_0 r^2} \mathbf{a}_R$$

which is the same as that of a point charge as one would expect.



## EXAMPLE

The finite sheet  $0 \leq x \leq 1$ ,  $0 \leq y \leq 1$  on the  $z = 0$  plane has a charge density  $\rho_S = xy(x^2 + y^2 + 25)^{3/2}$  nC/m<sup>2</sup>. Find

- (a) The total charge on the sheet
- (b) The electric field at  $(0, 0, 5)$
- (c) The force experienced by a  $-1$  mC charge located at  $(0, 0, 5)$

### **Solution:**

$$(a) \quad Q = \int \rho_S dS = \int_0^1 \int_0^1 xy(x^2 + y^2 + 25)^{3/2} dx dy \text{ nC}$$

Since  $x dx = 1/2 d(x^2)$ , we now integrate with respect to  $x^2$



## SOLUTION

$$\begin{aligned} Q &= \frac{1}{2} \int_0^1 y \int_0^1 (x^2 + y^2 + 25)^{3/2} d(x^2) dy \text{ nC} \\ &= \frac{1}{2} \int_0^1 y \frac{2}{5} (x^2 + y^2 + 25)^{5/2} \Big|_0^1 dy \\ &= \frac{1}{5} \int_0^1 \frac{1}{2} [(y^2 + 26)^{5/2} - (y^2 + 25)^{5/2}] d(y^2) \\ &= \frac{1}{10} \cdot \frac{2}{7} [(y^2 + 26)^{7/2} - (y^2 + 25)^{7/2}] \Big|_0^1 \\ &= \frac{1}{35} [(27)^{7/2} + (25)^{7/2} - 2(26)^{7/2}] \end{aligned}$$

$$Q = 33.15 \text{ nC}$$



## SOLUTION

$$(b) \mathbf{E} = \int \frac{\rho_S dS \mathbf{a}_R}{4\pi\epsilon_0 r^2} = \int \frac{\rho_S dS (\mathbf{r} - \mathbf{r}')}{4\pi\epsilon_0 |\mathbf{r} - \mathbf{r}'|^3}$$

where  $\mathbf{r} - \mathbf{r}' = (0, 0, 5) - (x, y, 0) = (-x, -y, 5)$ . Hence,

$$\begin{aligned} \mathbf{E} &= \int_0^1 \int_0^1 \frac{10^{-9} xy (x^2 + y^2 + 25)^{3/2} (-x\mathbf{a}_x - y\mathbf{a}_y + 5\mathbf{a}_z) dx dy}{4\pi \cdot \frac{10^{-9}}{36\pi} (x^2 + y^2 + 25)^{3/2}} \\ &= 9 \left[ -\int_0^1 x^2 dx \int_0^1 y dy \mathbf{a}_x - \int_0^1 x dx \int_0^1 y^2 dy \mathbf{a}_y + 5 \int_0^1 x dx \int_0^1 y dy \mathbf{a}_z \right] \\ &= 9 \left( \frac{-1}{6}, \frac{-1}{6}, \frac{5}{4} \right) \\ &= (-1.5, -1.5, 11.25) \text{ V/m} \end{aligned}$$

$$(c) \mathbf{F} = q\mathbf{E} = (1.5, 1.5, -11.25) \text{ mN}$$



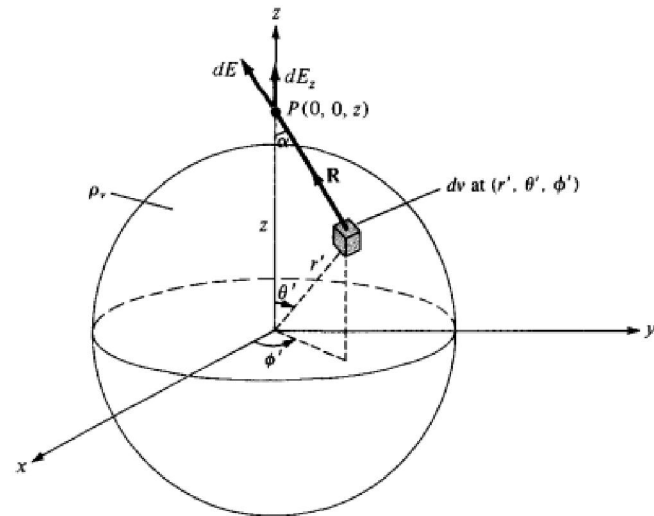
# VOLUME CHARGE

- Let the volume charge distribution with uniform charge density  $\rho_v$  be as shown in Figure
- The charge  $dQ$  associated with the elemental volume  $dv$  is

$$dQ = \rho_v dv$$

and hence the total charge in a sphere of radius  $a$  is

$$\begin{aligned} Q &= \int \rho_v dv = \rho_v \int dv \\ &= \rho_v \frac{4\pi a^3}{3} \end{aligned}$$



# VOLUME CHARGE

The electric field  $d\mathbf{E}$  at  $P(0, 0, z)$  due to the elementary volume charge is

$$d\mathbf{E} = \frac{\rho_v dv}{4\pi\epsilon_0 R^2} \mathbf{a}_R$$

where  $\mathbf{a}_R = \cos \alpha \mathbf{a}_z + \sin \alpha \mathbf{a}_\rho$ . Due to the symmetry of the charge distribution, the contributions to  $E_x$  or  $E_y$  add up to zero. We are left with only  $E_z$ , given by

$$E_z = \mathbf{E} \cdot \mathbf{a}_z = \int dE \cos \alpha = \frac{\rho_v}{4\pi\epsilon_0} \int \frac{dv \cos \alpha}{R^2}$$

Again, we need to derive expressions for  $dv$ ,  $R^2$ , and  $\cos \alpha$ .

$$dv = r'^2 \sin \theta' dr' d\theta' d\phi'$$

Applying the cosine rule to Figure we have

$$R^2 = z^2 + r'^2 - 2zr' \cos \theta'$$

$$r'^2 = z^2 + R^2 - 2zR \cos \alpha$$



# VOLUME CHARGE

$$\cos \alpha = \frac{z^2 + R^2 - r'^2}{2zR}$$

$$\cos \theta' = \frac{z^2 + r'^2 - R^2}{2zr'}$$

$$\sin \theta' d\theta' = \frac{R dR}{z r'}$$

Substituting

$$\begin{aligned} E_z &= \frac{\rho_v}{4\pi\epsilon_0} \int_{\phi'=0}^{2\pi} d\phi' \int_{r'=0}^a \int_{R=z-r'}^{z+r'} r'^2 \frac{R dR}{zr'} dr' \frac{z^2 + R^2 - r'^2}{2zR} \frac{1}{R^2} \\ &= \frac{\rho_v 2\pi}{8\pi\epsilon_0 z^2} \int_{r'=0}^a \int_{R=z-r'}^{z+r'} r' \left[ 1 + \frac{z^2 - r'^2}{R^2} \right] dR dr' \\ &= \frac{\rho_v \pi}{4\pi\epsilon_0 z^2} \int_0^a r' \left[ R - \frac{(z^2 - r'^2)}{R} \right]_{z-r'}^{z+r'} dr' \\ &= \frac{\rho_v \pi}{4\pi\epsilon_0 z^2} \int_0^a 4r'^2 dr' = \frac{1}{4\pi\epsilon_0} \frac{1}{z^2} \left( \frac{4}{3} \pi a^3 \rho_v \right) \end{aligned}$$

or

$$\mathbf{E} = \frac{Q}{4\pi\epsilon_0 z^2} \mathbf{a}_z$$

